

Spin memristive systems: Spin memory effects in semiconductor spintronics

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Recently, in addition to the well-known resistor, capacitor, and inductor, a fourth passive circuit element, named memristor, has been identified following theoretical predictions. The model example used in such case consisted in a nanoscale system with coupled ionic and electronic transport. Here, we discuss a system whose memristive behavior is based entirely on the electron-spin degree of freedom, which allows for a more convenient control than the ionic transport in nanostructures. An analysis of time-dependent spin transport at a semiconductor/ferromagnet junction provides a direct evidence of memristive behavior. Our scheme is fundamentally different from previously discussed schemes of memristive systems and broadens the possible range of applications of semiconductor spintronics.

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In 1971 Leon Chua,¹ analyzing mathematical relations between pairs of fundamental circuit variables, proposed a fourth two-terminal circuit element characterized by a relationship between the charge and the flux. He called that element a memristor (or memory resistor). Five years later, Chua and Kang² introduced a more general class of two-terminal devices: memristive systems. If w denotes a set of n -state variables describing the internal state of the system, an n th-order current-controlled memristive system is described by the equations

$$V = R(w, i, t)i, \quad (1)$$

$$\dot{w} = f(w, i, t), \quad (2)$$

where V and i denote the voltage and current and R is a generalized resistance. The equation for a memristor is a particular case of Eqs. (1) and (2), when R depends only on charge, namely,

$$V = R[q(t)]i. \quad (3)$$

Although several experimental systems were suggested to behave as memristive systems,² the real interest in these devices resurfaced only recently when a group of scientists from Hewlett Packard identified a specific experimental realization of memristor.³ In their scheme, a memory effect is achieved in solid-state thin film two-terminal devices. In that case, the memristive behavior is based on the coupling between transport of atomic degrees of freedom (e.g., oxygen vacancies acting as mobile dopants defining the internal state of the device) and of electrons. The electron current flowing through such a device dynamically changes the internal state of the latter, which, in turn, influences the electron transport in a nonlinear way.

Ionic transport, however, relies on microscopic features of the system which are generally difficult to control, especially at the nanoscale. The use of a fundamentally different degree of freedom, which allows for the realization of memristive behavior, would be thus desirable.

In this work we demonstrate that such degree of freedom is provided by the electron spin and show that memristive

behavior is common for the broad class of semiconductor spintronic devices. This class involves systems whose transport properties depend on the level of electron-spin polarization in a semiconductor, which is influenced by an external control parameter (such as an applied voltage). The adjustment in electron-spin polarization to a variation in the control parameter takes some time (this process involves diffusion and/or relaxation of electron-spin polarization), resulting in a memory effect accompanied by all other requirements for memristive behavior, such as absence of energy storage.

As an example of semiconductor spintronic system exhib-

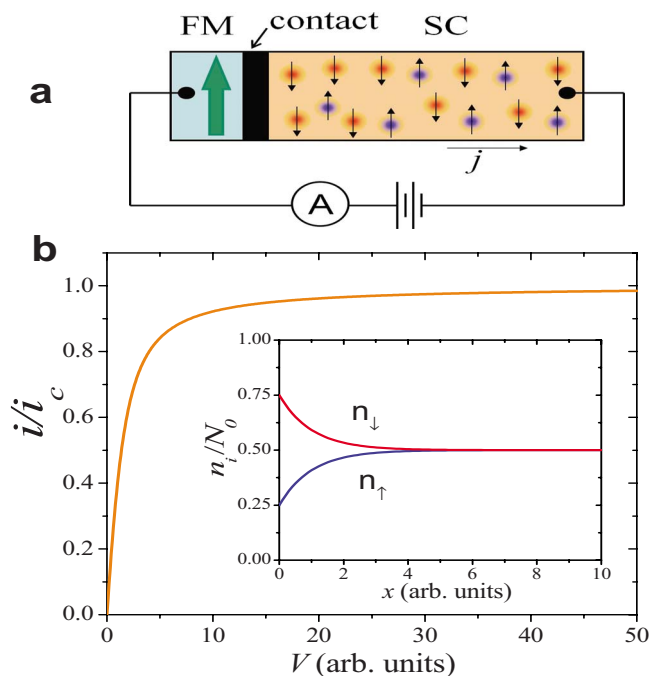


FIG. 1. (Color online) Semiconductor/half-metal junction. (a) Schematic representation of the circuit made of an interface between a semiconductor and a half metal. (b) Typical dc current-voltage characteristics. Inset: spin-up and spin-down densities in the semiconductor region as a function of the distance from the contact. $i_c = eN_0\sqrt{D}/(2\tau_s)$ is the critical current density (Ref. 5).

iting memristive behavior, let us consider a semiconductor/half-metal junction as schematically shown in Fig. 1(a). We consider a junction with half metals (ferromagnets with 100% spin polarization at the Fermi level) because these act as perfect spin filters and, therefore, are more sensitive to the level of electron-spin polarization. However, we also expect that the prediction of memristive behavior is valid for junctions with ferromagnets having less than 100% spin polarization as a result of anticipated spin-based peculiarity in the i - V curve of these systems.⁴ Understanding the properties of the systems discussed in this Brief Report is of great interest in the context of actively studied spin-injection/spin-extraction processes at semiconductor/ferromagnet junctions.⁴⁻¹⁴

We will be interested mainly in the process of spin extraction, where the electron flow is from the semiconductor into the half metal, which is especially interesting because of the recently predicted spin-blockade phenomenon in such junctions.^{4,5} The physics of the spin blockade is the following: the half metal accepts electrons of only one—say up—spin direction. Spin-down electrons cannot enter the half metal and, therefore, form a cloud near the contact [see the inset in Fig. 1(b)] when a current flows through the system. This cloud increases with increasing current. At a critical current density the density of spin-up electrons near the contact becomes insufficient to provide a further current increase. In other words, transport of spin-up electrons through the contact becomes blocked by the cloud of spin-down electrons near the contact. It was predicted by the present authors⁴ that the spin blockade leads to a saturated i - V curve

as that shown in Fig. 1(b). Here, we show that such a system has all the necessary components to exhibit memristive behavior.

In order to simulate current-voltage characteristics of the junction we assume that the applied voltage drops mainly on the semiconductor part and contact region. Taking a constant conductivity of the semiconductor region and contact conductivity proportional to spin-up electron density near the contact,⁴ we can write

$$V = V_s + V_c = \left(\rho_s L + \rho_c^0 \frac{N_0}{2n_{\uparrow}(0)} \right) i, \quad (4)$$

where ρ_s and L are the semiconductor resistivity and length, ρ_c^0 is the contact resistivity at $V \rightarrow 0$, N_0 is the electron density in the semiconductor, and $n_{\uparrow}(0)$ is the density of spin-up electrons near the contact. For simplicity, we assume constant electron density in the semiconductor, i.e., $n_{\uparrow} + n_{\downarrow} = N_0$.

The electron-spin densities in the semiconductor are described by the spin drift-diffusion model¹⁴

$$e \frac{\partial n_{\uparrow(\downarrow)}}{\partial t} = \text{div } \vec{i}_{\uparrow(\downarrow)} + \frac{e}{2\tau_{sf}} (n_{\downarrow(\uparrow)} - n_{\uparrow(\downarrow)}), \quad (5)$$

$$\vec{i}_{\uparrow(\downarrow)} = \sigma_{\uparrow(\downarrow)} \vec{E} + eD \nabla n_{\uparrow(\downarrow)}. \quad (6)$$

Here, $\vec{i}_{\uparrow(\downarrow)}$ is the current density of spin-up (-down) electrons, τ_{sf} is the spin-relaxation time, $\sigma_{\uparrow(\downarrow)}$ is the conductivity of spin-up (-down) electrons, and D is the diffusion coefficient.

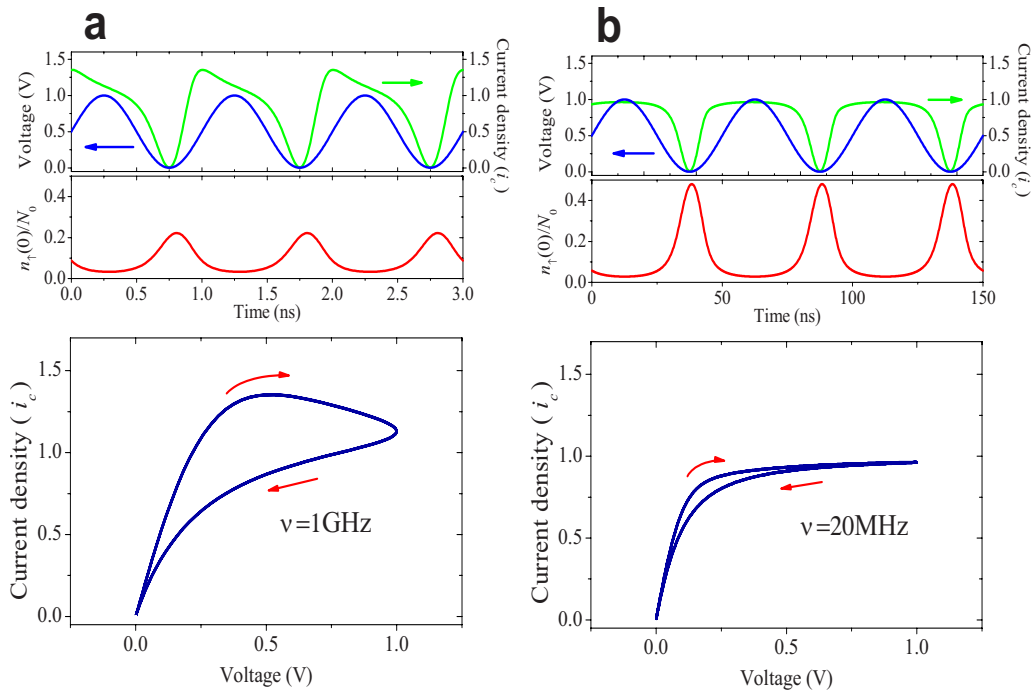


FIG. 2. (Color online) Simulations of ac response of the system. The applied voltages (blue lines) are $V = V_1 + V_2 \sin(2\pi\nu t)$ with $V_1 = V_2 = 0.5$ V, $\nu = 10^9$ Hz in (a) and $\nu = 2 \cdot 10^7$ Hz in (b). In both cases the current densities (green lines in units of the critical current density) and spin-up electron densities near the contact (red lines) show saturation typical for spin blockade (Ref. 5). It is clearly seen that i - V hysteresis is significantly reduced in the low-frequency case. The calculations were made using the following system parameters: $L = 20$ μm , $D = 220$ cm^2/s , $\mu = 8500$ $\text{cm}^2/(\text{V}\cdot\text{s})$, $\tau_{sf} = 2$ ns, $N_0 = 5 \times 10^{15}$ cm^{-3} , and $\rho_c^0/(\rho_s L) = 1$.

According to the present classification scheme for memristive behavior,² it is clear that the semiconductor/half-metal junction is a current-controlled memristive system. This can be seen by comparing directly Eqs. (4) and (5) with Eqs. (1) and (2), respectively [we can certainly write Eq. (5) in terms of n discrete variables]. In order to observe memristive behavior, we solved Eqs. (4)–(6) self-consistently at every given time¹⁵ with the following boundary conditions: $i_{\uparrow}(0) = i$, $i_{\downarrow}(0) = 0$, and $i_{\uparrow}(L) = i_{\downarrow}(L) = i/2$ corresponding to the process of spin extraction at the contact.

Figure 2 shows results of our simulations with a time-dependent applied voltage $V = V_1 + V_2 \sin(2\pi\nu t)$ of high (1 GHz) and low (20 MHz) frequency ν with $V_1 \geq V_2$. Such voltage profile has been selected in order to be always in the spin-extraction regime. The calculated i - V curves (bottom panels in Fig. 2) exhibit a frequency-dependent hysteretic behavior typical of memristive systems. In particular, we can readily notice the distinctive zero-crossing property of these curves: no current flows through the structure when the voltage drop is zero (i - V curves pass through the $i=0$, $V=0$ point). This property is related to the fact that there is no energy storage in our system as opposed to the energy storage in the usual capacitive or inductive circuit elements.

Another interesting feature of the i - V curves is their frequency behavior. It follows from Fig. 2 that the hysteresis is significantly suppressed at 20 MHz frequency. This is a manifestation of the fact that at low frequencies our system behaves essentially as a nonlinear resistor. Physically, at low applied voltage frequencies, the electron-spin polarization in the semiconductor has enough time to adjust to any present value of the voltage. Therefore, the current through the system at low frequencies is nonlinear (because of spin blockade) but essentially history independent (no hysteresis). It is also clear that at very high frequencies ($\nu \rightarrow \infty$) the electron-spin polarization does not have any time for redistribution within the oscillation period. Therefore, the spin-up density near the contact and, correspondingly, the contact resistivity can be considered constant, so that the device operates as a linear resistor. The above described frequency behavior of the semiconductor/half-metal junction is typical for memristive systems.²

Although the model represented by Eqs. (4)–(6) takes into account the main physics of the underlying system operation, it is still quite complex for circuit analysis. Here, we thus discuss a simplified (and more transparent) model which can be obtained in the following way. Instead of tracking the whole spin-density distribution in the semiconductor, let us focus our attention on the integrated spin density (surface spin density) $N_{\uparrow(\downarrow)} = \int_0^L n_{\uparrow(\downarrow)} dx$ and select this quantity as the only system parameter. Integrating Eq. (5) from 0 to L (i.e., over the whole semiconductor length) we obtain the following equation:

$$e \frac{\partial N_{\uparrow}}{\partial t} = -\frac{1}{2}i + \frac{e}{\tau_{sf}} \left(\frac{NL}{2} - N_{\uparrow} \right). \quad (7)$$

Equation (7) simply states that the change in the integrated spin-up density is due to the injection/extraction of spin-up electrons through the semiconductor boundaries [the

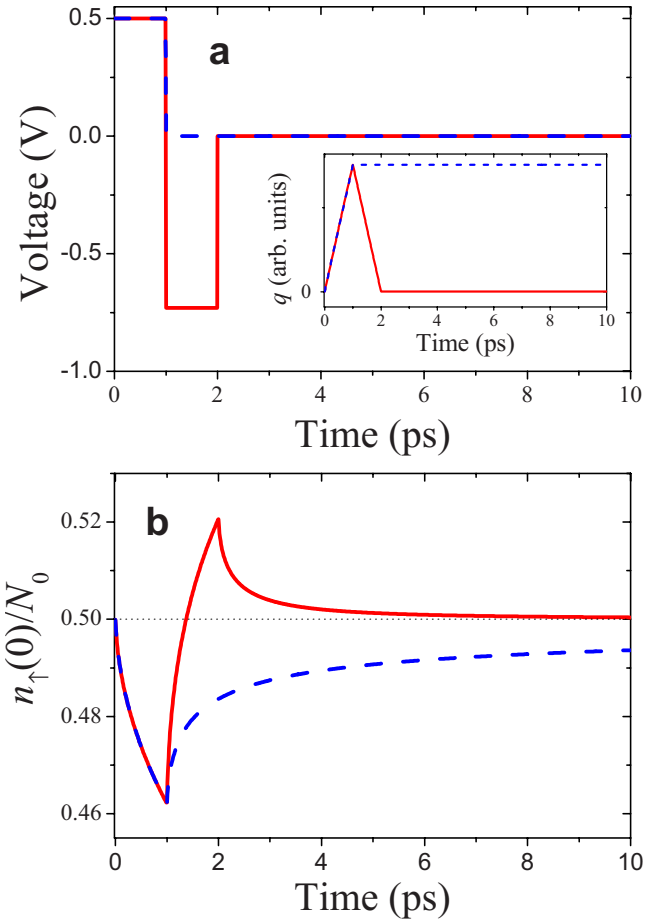


FIG. 3. (Color online) System dynamics excited by step voltages. (a) Unipolar (blue line) and bipolar (red line) step voltage profiles used in our calculations. The spin-injection process (at negative applied voltage) was modeled using a constant interface resistance model. Inset: Total charge flown through the system as a function of time. The profile of bipolar voltage was selected in such a way that the corresponding total charge $q=0$ for $t > 2$ ps (see the red line). (b) Evolution of spin-up electron density near the interface. In the case of bipolar voltage excitation (red line), the final value of spin-up density is close to its initial value (shown by the dotted line). This is a manifestation of a nearly perfect memristor behavior.

first “current” term in the right-hand side (RHS) of Eq. (7)] and spin-relaxation processes [the second term in the RHS of Eq. (7)]. Although the knowledge of N_{\uparrow} is not sufficient to exactly obtain $n_{\uparrow}(0)$, which enters into Eq. (4), we can approximately write

$$n_{\uparrow}(0) = g(N_{\uparrow}), \quad (8)$$

where g is a given smooth function. The simplified circuit model given by Eqs. (4), (7), and (8) describes again a memristive system. However, at short times, when spin-relaxation processes are not important the last term in the RHS of Eq. (7) can be neglected. The resulting set of equations describes then a perfect *memristor*. Indeed, integrating Eq. (7) we obtain

$$N_{\uparrow}(t) = N_{\uparrow}(0) - \frac{1}{2e} \int_0^t i(\tau) d\tau = N_{\uparrow}(0) - \frac{1}{2e} q(t), \quad (9)$$

and, substituting Eqs. (8) and (9) into Eq. (4), we get

$$V = \left[\rho_s L + \rho_c^0 \frac{N_0}{2g(N_{\uparrow}(0) - \frac{1}{2e} q(t))} \right]. \quad (10)$$

It follows from Eq. (10), which is a particular case of Eq. (3), that the resistivity of the system depends only on the amount of charge $q(t)$ flown through it.

In order to test the simplified circuit model predictions, we calculate, using Eqs. (4)–(6), the short-time system response to unipolar and bipolar voltage step excitations. Figure 3(b) demonstrates that our structure behaves almost as a perfect memristor at short times. In particular, the red line in Fig. 3(b) shows that when the total charge flown through the structure is equal to zero, the value of the system parameter $n_{\uparrow}(0)$ (and consequently of the device resistivity) becomes very close to its initial value already at $t=5$ ps. The opposite behavior is demonstrated by the blue line in Fig. 3(b): when the total charge flowing in the system is different from zero, the value of the spin-up density at the interface $n_{\uparrow}(0)$ is quite different from its initial value even at $t=10$ ps.

Even though the focus of this Brief Report is on the physical mechanisms underlying the spin memristive behavior, we briefly discuss possible applications of these results. Certainly, since the spin-relaxation time in semiconductors is quite short, the use of spin memristive systems in memory applications is unlikely. The problem is that the refreshing

rate for spin memristive systems should be much higher than those in usual dynamic random access memory (64 ms) unless a system with much longer spin lifetime is used. Instead, a more promising application area is analog electronics, in which utilization of spin memristive systems may be beneficial because of their unique characteristics. Finally, another interesting direction is in solid-state modeling and realization of neural networks and mechanisms of learning.

To summarize, we have demonstrated that a semiconductor/half-metal junction is in fact a memristive system. The origin of its unusual behavior in a circuit is completely based on the electronic spin degree of freedom, which is much easier to control than ionic transport at the nanoscale. In our scheme, the principal role is played by electron-spin diffusion and relaxation processes which drive the system to equilibrium. Moreover, we would like to note that the spin memristive behavior is not limited to the discussed structure; it should be typical of many semiconductor spintronic devices that involve spin filters, general semiconductor-ferromagnet junctions, etc. We also expect the memristive behavior to be more pronounced in structures with low electron density, where the level of electron-spin polarization can be significantly varied by external control parameters. We thus believe that our demonstration of memristive effects in spintronics is an important step forward in future practical applications of the newly discovered fourth circuit element.

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